

UCLA 2019

SANTA FE JETS AND HEAVY FLAVOR WORKSHOP

January 28-30, 2019

8:30am-5:30pm

UCLA IDRE Portal

Room 5628 Math Sciences Building

Hadrons' Partonic Structure from ab initio Lattice QCD Calculations

A status report

Jianwei Qiu

Theory Center, Jefferson Lab

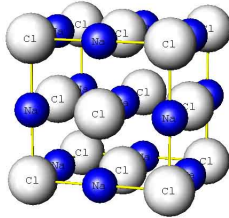
Theory Center

Jefferson Lab
EXPLORING THE NATURE OF MATTER

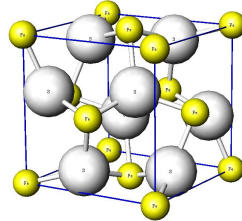
Hadron's partonic structure in QCD

❑ Structure – “a still picture”

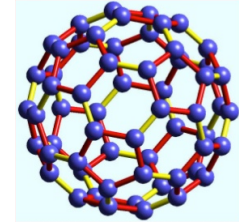
Crystal
Structure:



NaCl,
B1 type structure FeS₂,
C2, pyrite type structure



Nano-
material:

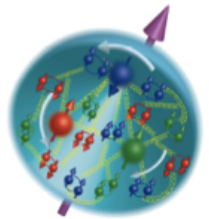


Fullerene, C₆₀

Motion of nuclei is much slower than the speed of light!

❑ No “still picture” for hadron's partonic structure!

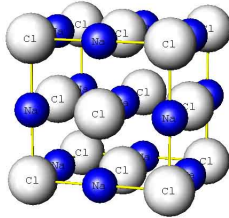
Motion of quarks/gluons is relativistic, color is fully entangled!



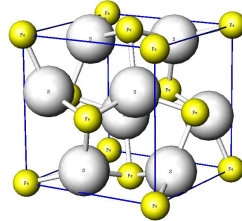
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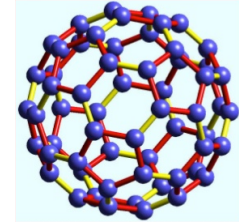


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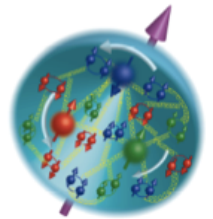


Fullerene, C60

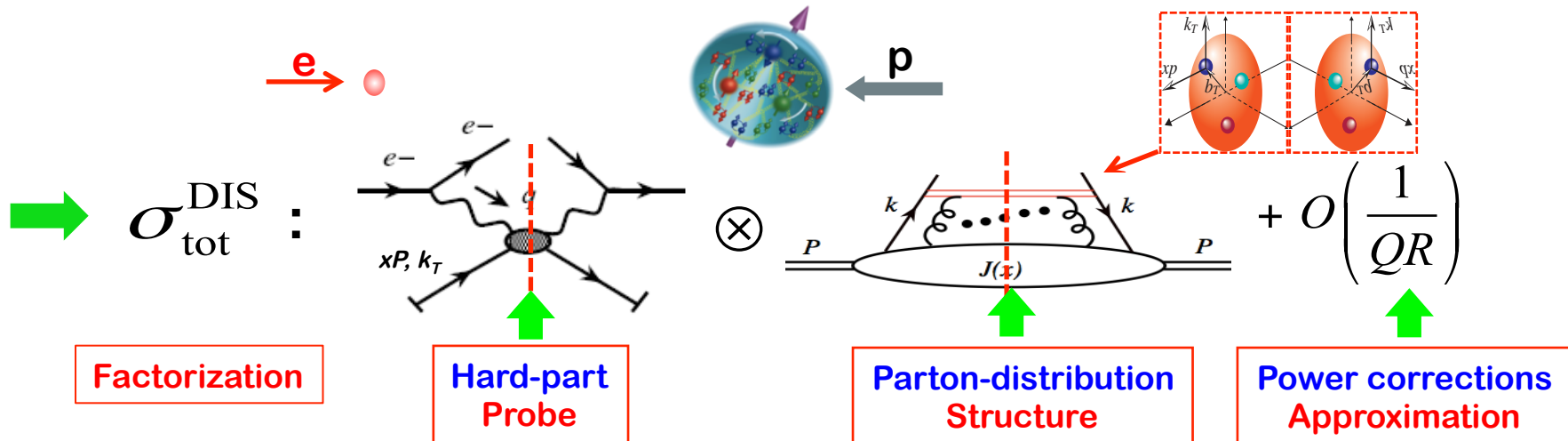
Motion of nuclei is much slower than the speed of light!

❑ No “still picture” for hadron's partonic structure!

Motion of quarks/gluons is relativistic, color is fully entangled!



❑ Asymptotic Freedom → Factorization → localized “probe”



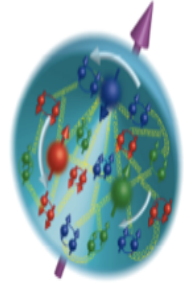
Hadron's partonic structure in QCD

□ Quantifying the partonic structure:

= Matrix elements of quarks and/or gluons

= Quantum “probabilities” $\langle P, S | \mathcal{O}(\bar{\psi}, \psi, A^\mu) | P, S \rangle$

None of these matrix elements is a direct physical observable in QCD – color confinement!



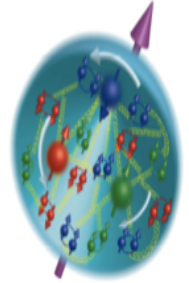
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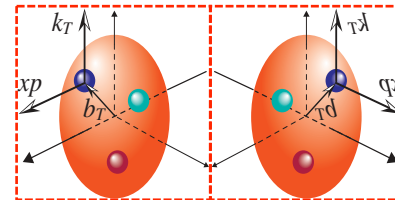
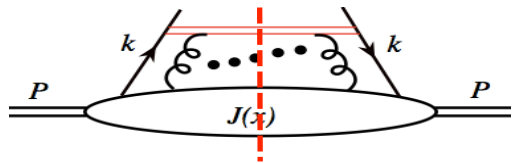
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Parton distribution functions (PDFs):



$$f_q(x, \mu^2) \equiv \int \frac{dP^+ \xi^-}{2\pi} e^{-ixP^+ \xi^-} \langle P | \bar{\psi}(\xi^-) \frac{\gamma^+}{2P^+} \exp \left\{ -ig \int_0^{\xi^-} d\eta^- A^+(\eta^-) \right\} \psi(0) | P \rangle$$

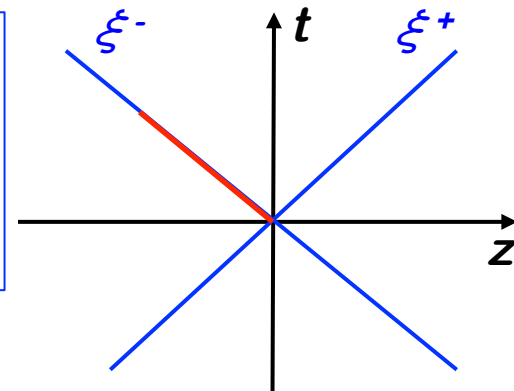
Dominated by the region:

$$\xi^- \lesssim 1/(xP^+) \sim 1/Q$$

Interpreted as:

*Probability density
to find a quark with a momentum fraction x*

Quantum correlation
of quark fields
along ξ^- direction!
(Conjugated to
the large P^+)



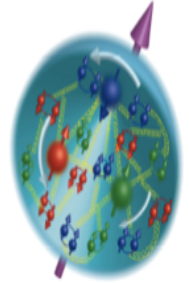
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Multi-parton quantum correlations:

✧ Spin-dependent cross section (one hard scale):

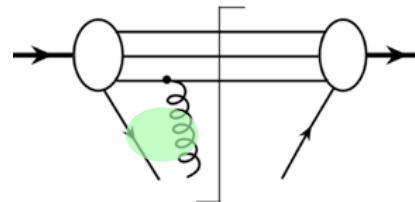
$$\sigma(Q, \vec{s}) \propto \left| \underbrace{\text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} + \dots}_{\text{Quantum interference}} \right|^2$$

The diagrams show various partonic interactions. The first diagram has a quark line with momentum k and a gluon line with momentum $t \sim 1/Q$. The second diagram shows a different partonic configuration. The third diagram shows a more complex interaction involving multiple partons.

Quantum interference

✧ Spin-Asymmetry:

$$\sigma(s) - \sigma(-s) \quad \rightarrow \quad T^{(3)}(x, x) \propto$$



Quantum interference between a quark state and a quark/gluon composite state – Twist-3 operators!

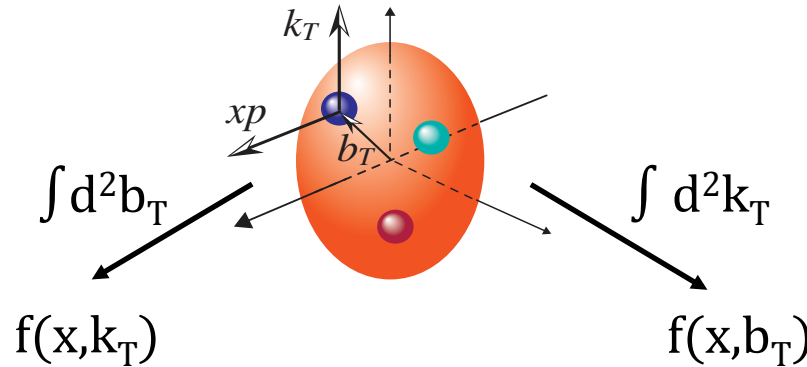
Hadron's partonic structure in QCD

□ Beyond collinear PDFs – 3D confined motion and spatial imaging:

*Momentum
Space*

TMDs

*Confined
motion*



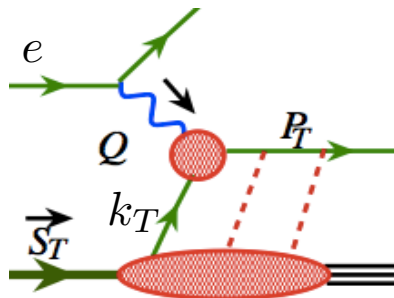
*Coordinate
Space*

GPDs

*Confined
spatial
distribution*

Two-scale observables!

✧ Semi-inclusive process (SIDIS):



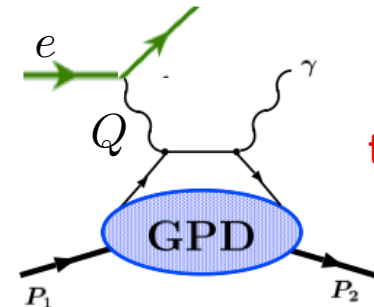
*Observe a
final-state
Hadron
(Control P_T)*

Two scales: $Q \gg P_T$

Initial-state proton is broken!

Probe “ P_T ” gives the “ k_T ”

✧ Exclusive process (DVCS):



$t = (p_1 - p_2)^2$

Two scales: $Q^2 \gg |t|^2$

Initial-state proton is NOT broken!

F.T. of “ t ” gives the “ b_T ”

Global QCD analyses – A successful story

□ World data with “Q” > 2 GeV

+ QCD Factorization:

e-H: $F_2(x_B, Q^2) = \sum_f C_f(x_B/x, \mu^2/Q^2) \otimes f(x, \mu^2)$

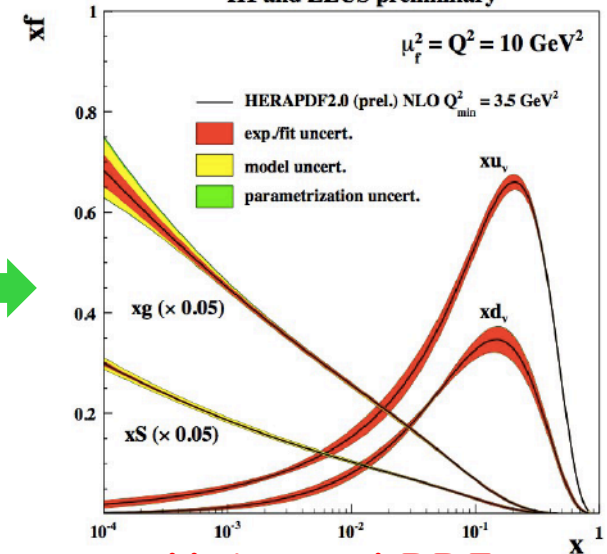
H-H: $\frac{d\sigma}{dy dp_T^2} = \sum_{ff'} f(x) \otimes \frac{d\hat{\sigma}_{ff'}}{dy dp_T^2} \otimes f'(x')$



+ DGLAP Evolution:

$$\frac{\partial f(x, \mu^2)}{\partial \ln \mu^2} = \sum_{f'} P_{ff'}(x/x') \otimes f'(x', \mu^2)$$

See also Petriello's talk
H1 and ZEUS preliminary




Universal PDFs

Global QCD analyses – A successful story

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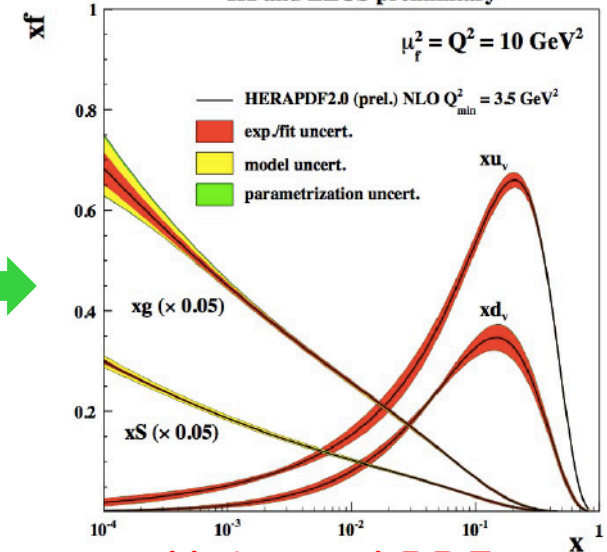
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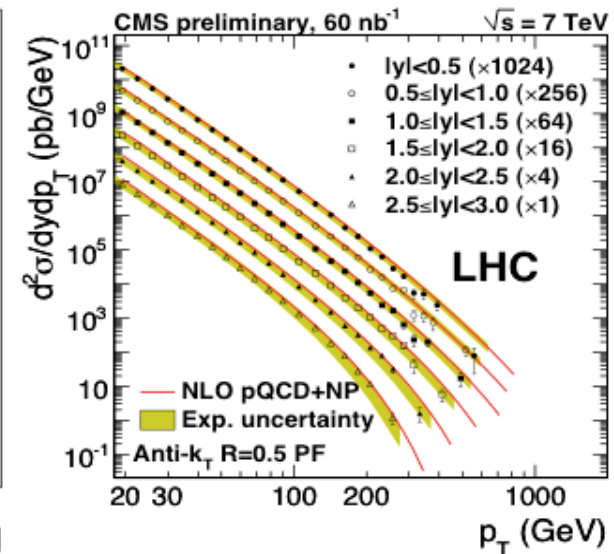
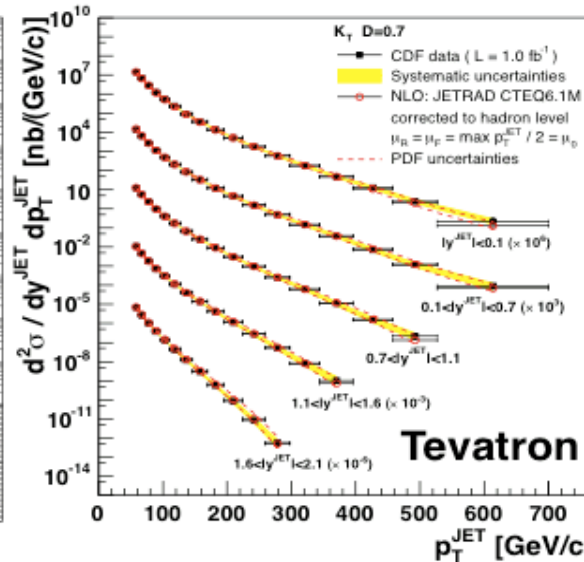
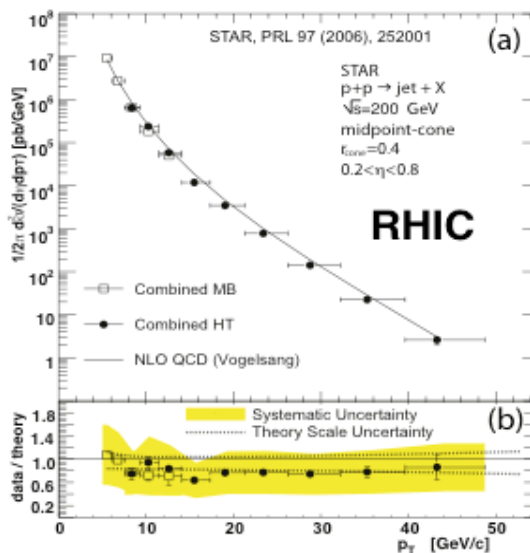
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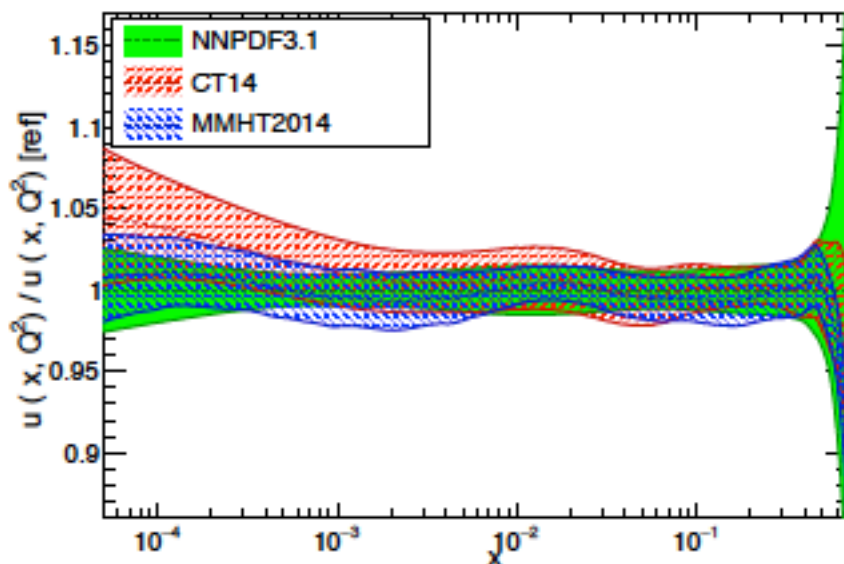
Universal PDFs



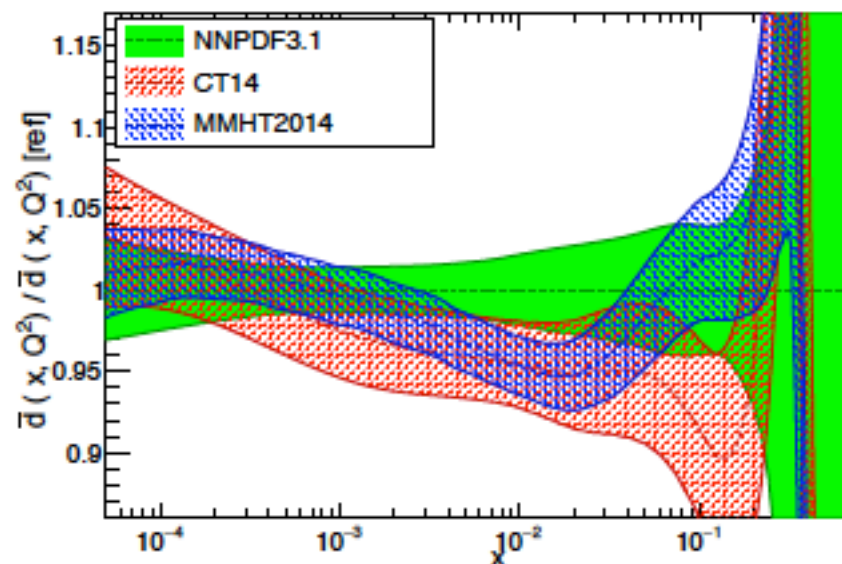
Uncertainties of PDFs

1711.07916

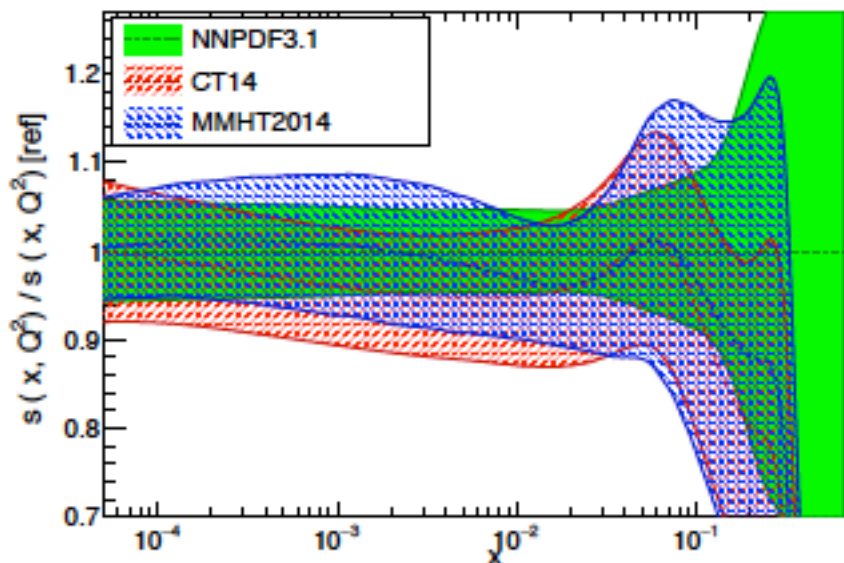
NNLO, $Q = 100$ GeV



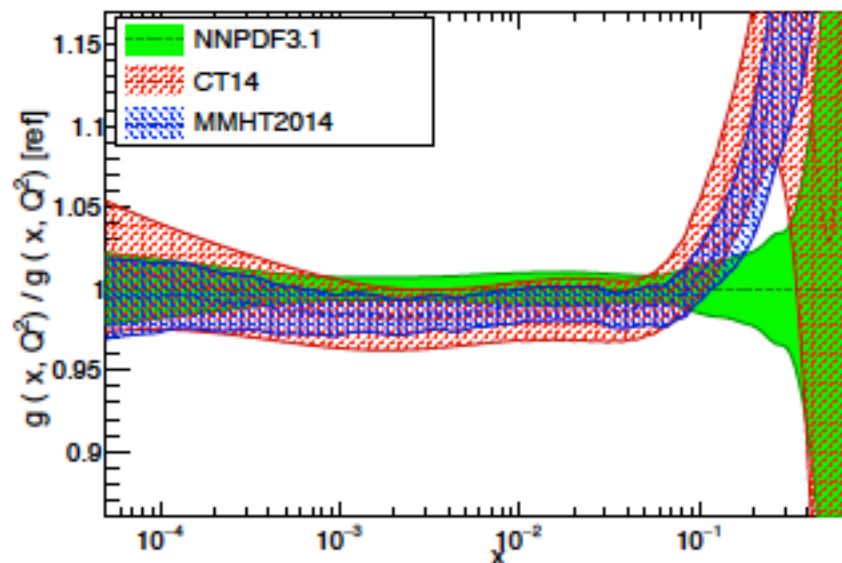
NNLO, $Q = 100$ GeV



NNLO, $Q = 100$ GeV

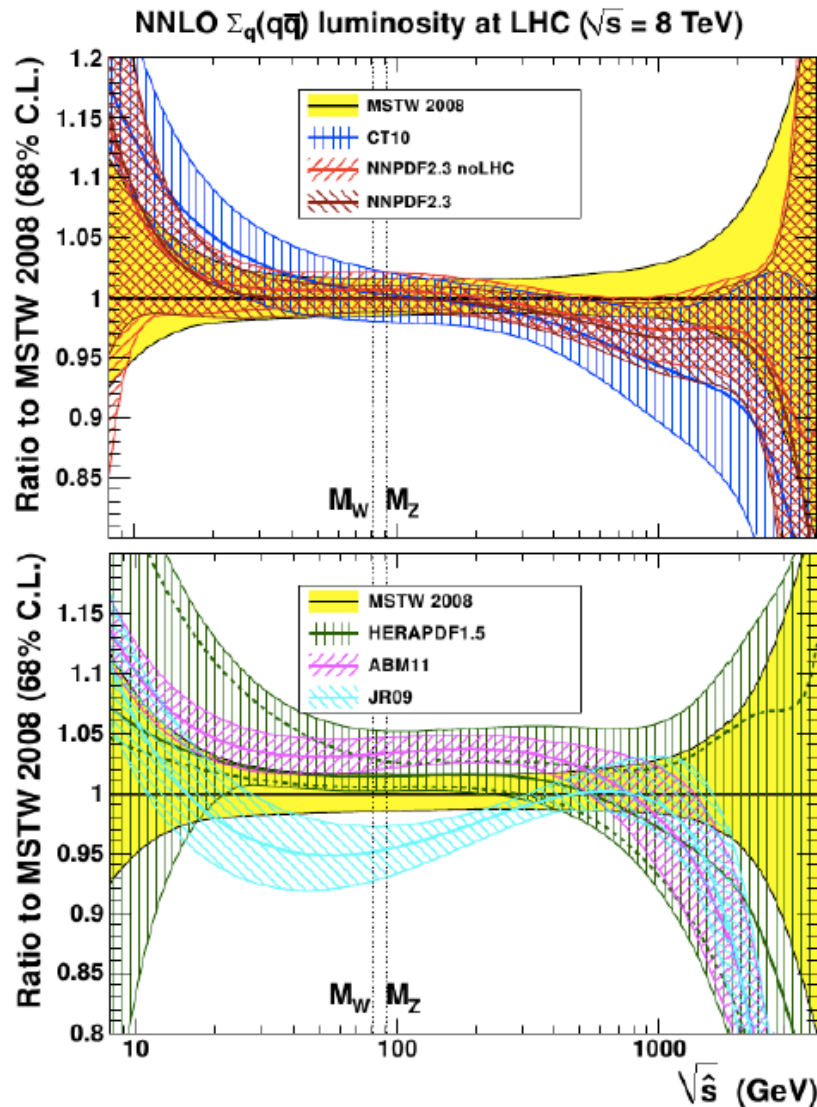


NNLO, $Q = 100$ GeV

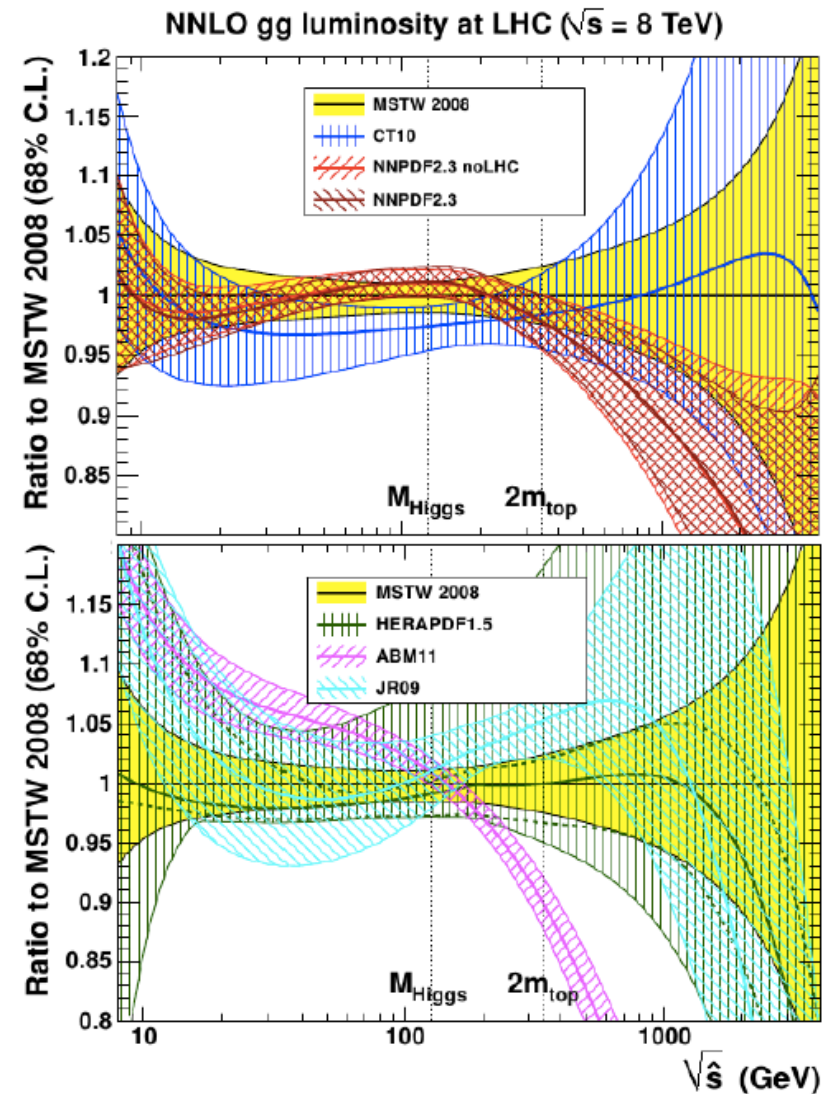


Partonic luminosities – discovery potential

$q - q\bar{q}$

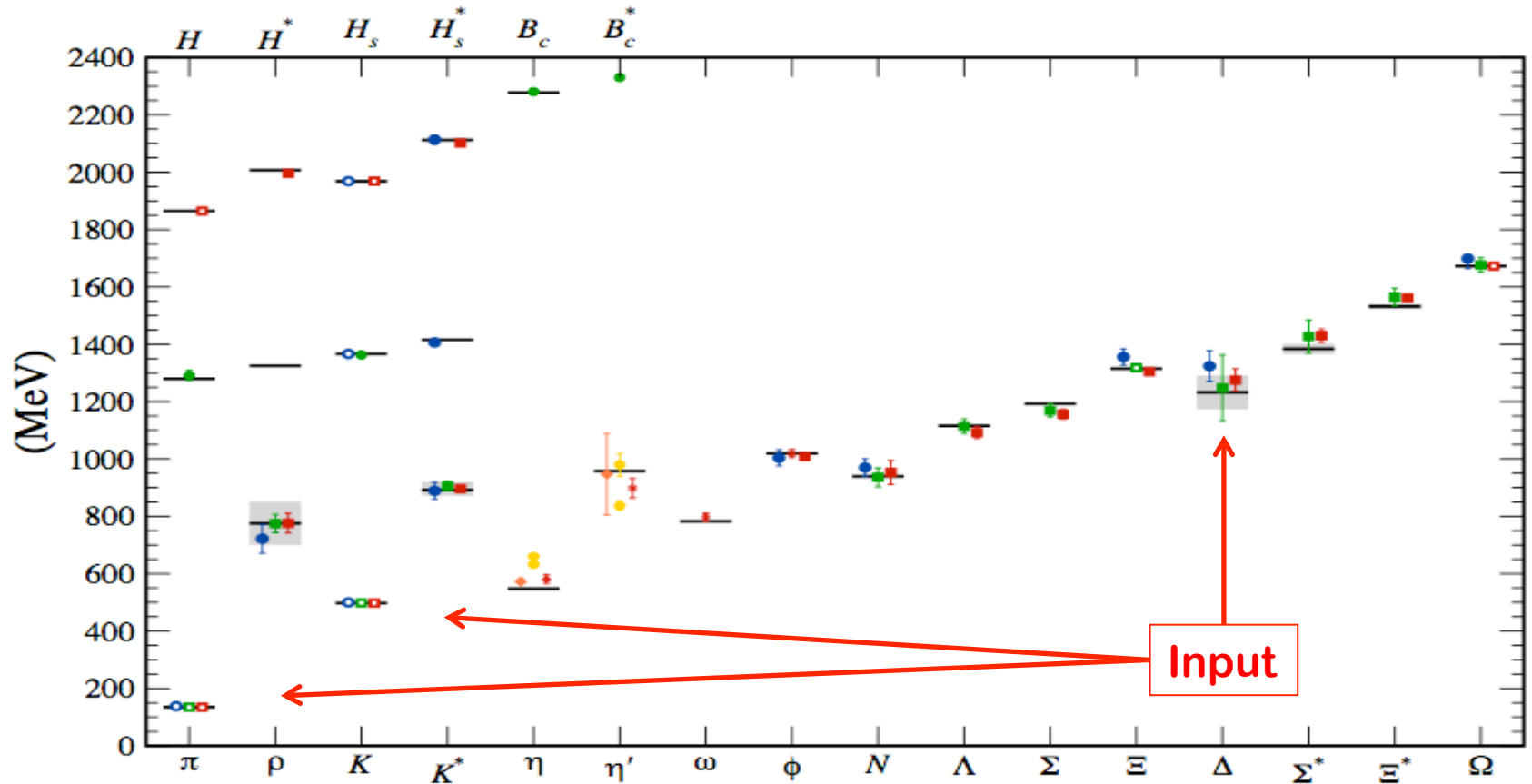


See also Petriello's talk
 $g - g$



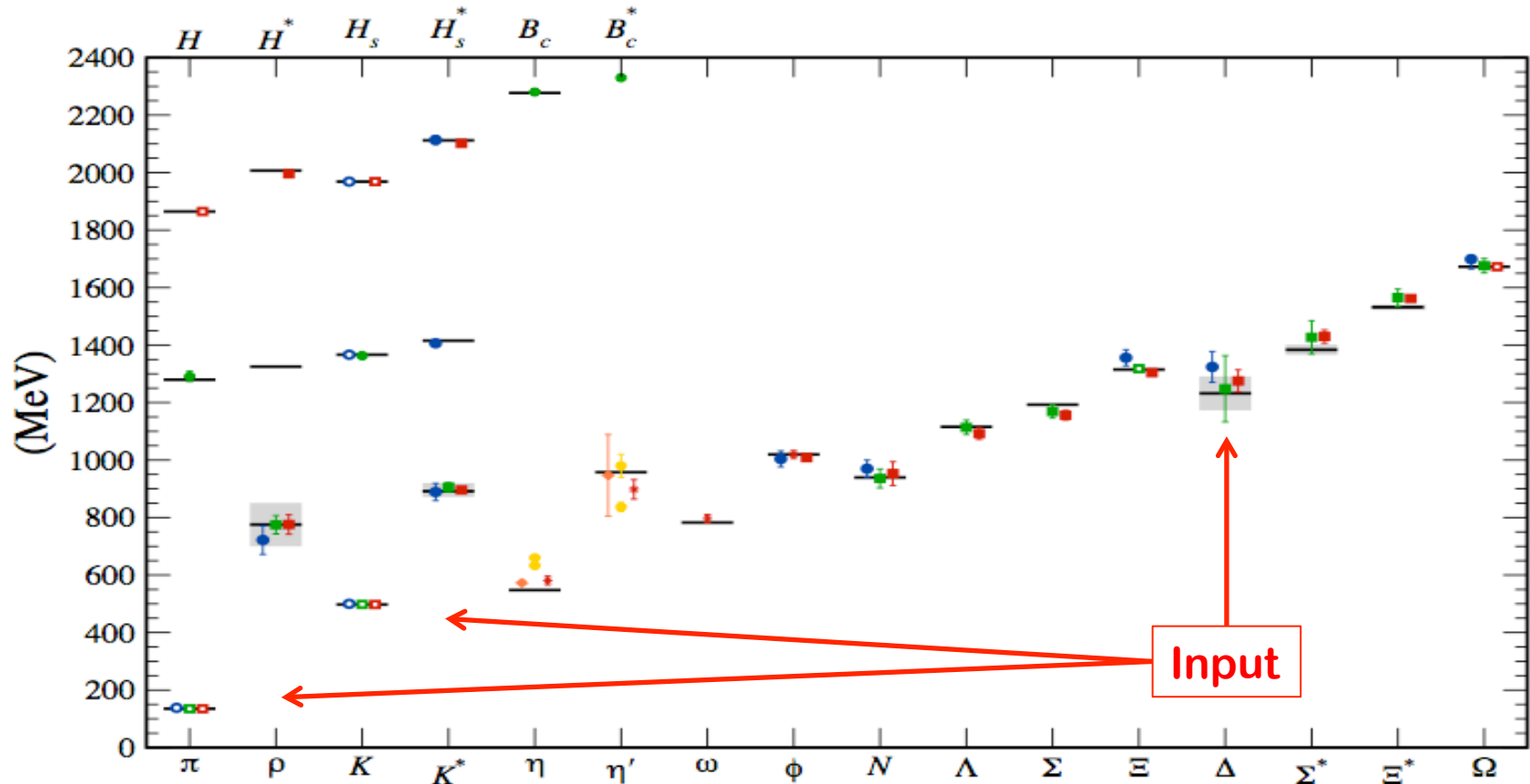
Lattice QCD

□ Hadron masses: Predictions with limited inputs



Lattice QCD

- Hadron masses: Predictions with limited inputs



- Lattice “time” is Euclidean: $\tau = i t$

Cannot calculate PDFs, TMDs, ..., directly, whose operators are time-dependent!

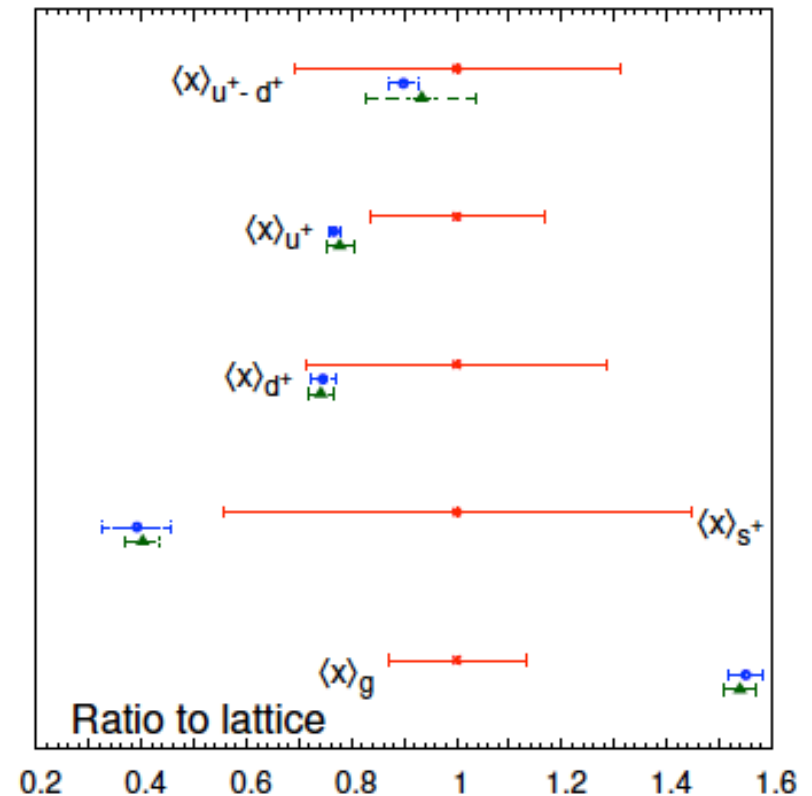
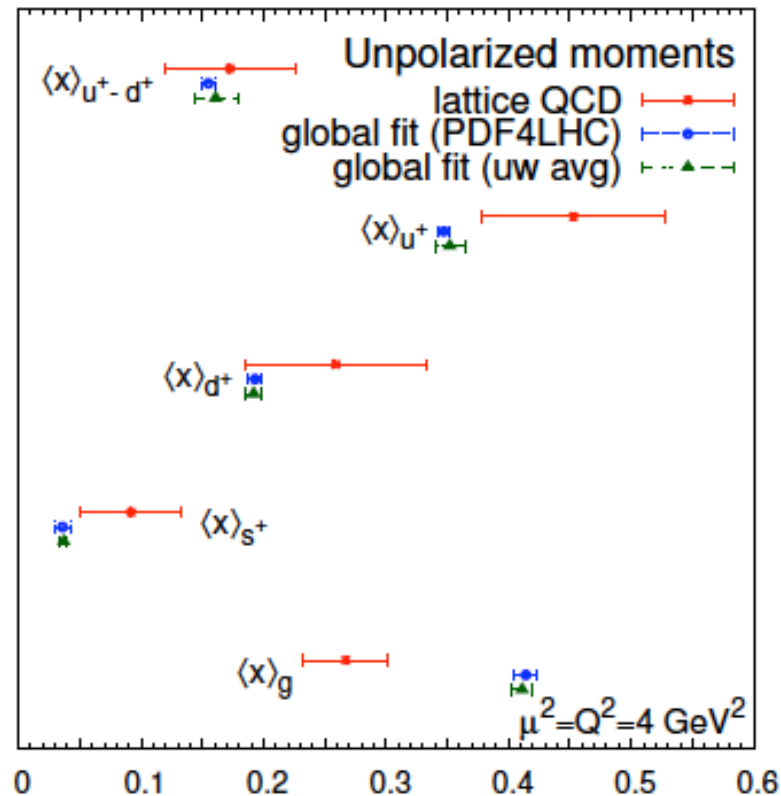
Moments of PDFs from lattice QCD

□ Moments of PDFs – matrix elements of local operators

$$\langle x^n(\mu^2) \rangle_q \equiv \int_0^1 dx x^n q(x, \mu^2)$$

$$q^\pm \equiv q \pm \bar{q} \quad \text{and} \quad \Delta q^\pm \equiv \Delta q \pm \Delta \bar{q}$$

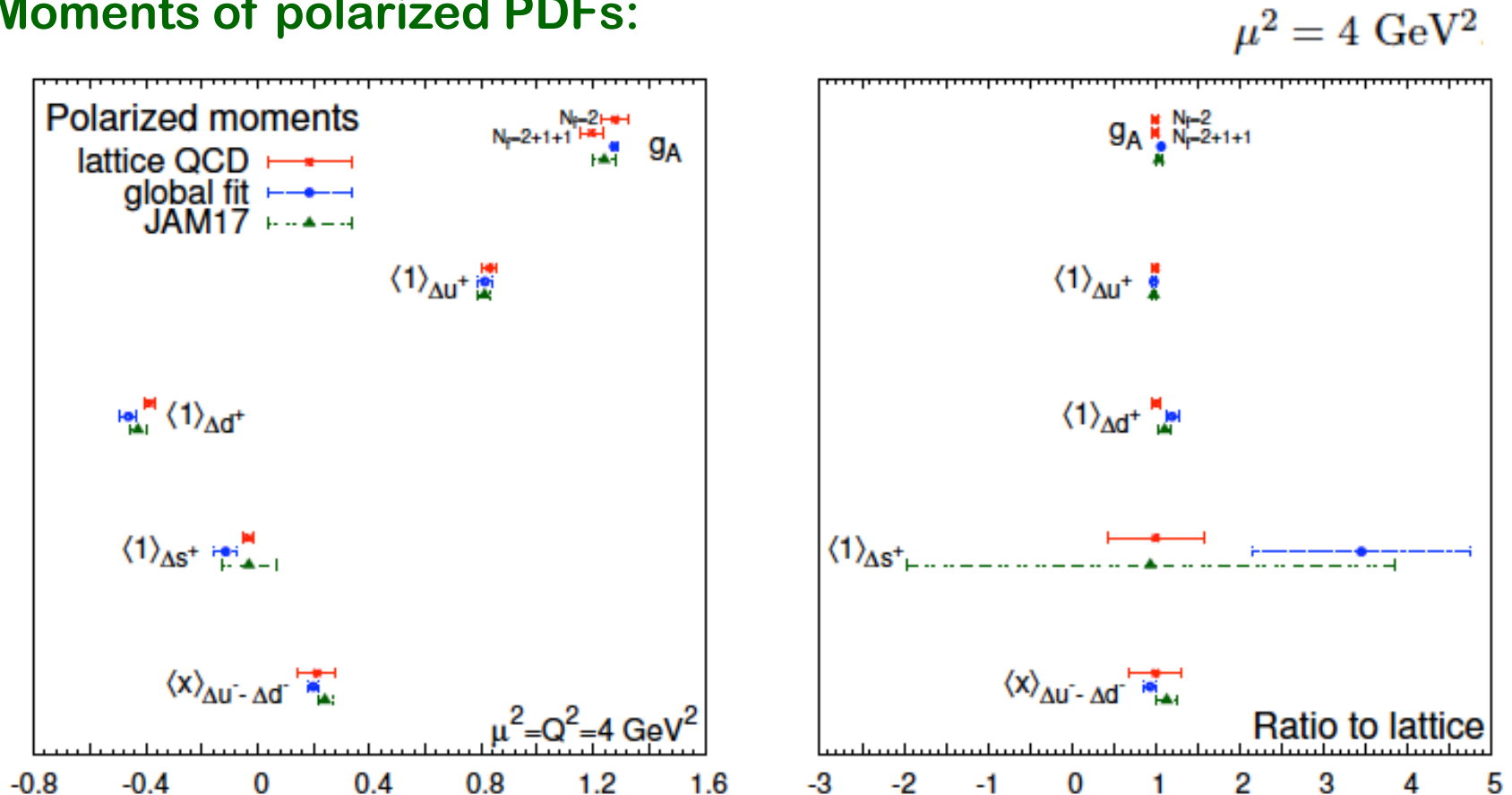
□ Moments of unpolarized PDFs:



$\mu^2 = 4 \text{ GeV}^2$

Moments of PDFs from lattice QCD

□ Moments of polarized PDFs:



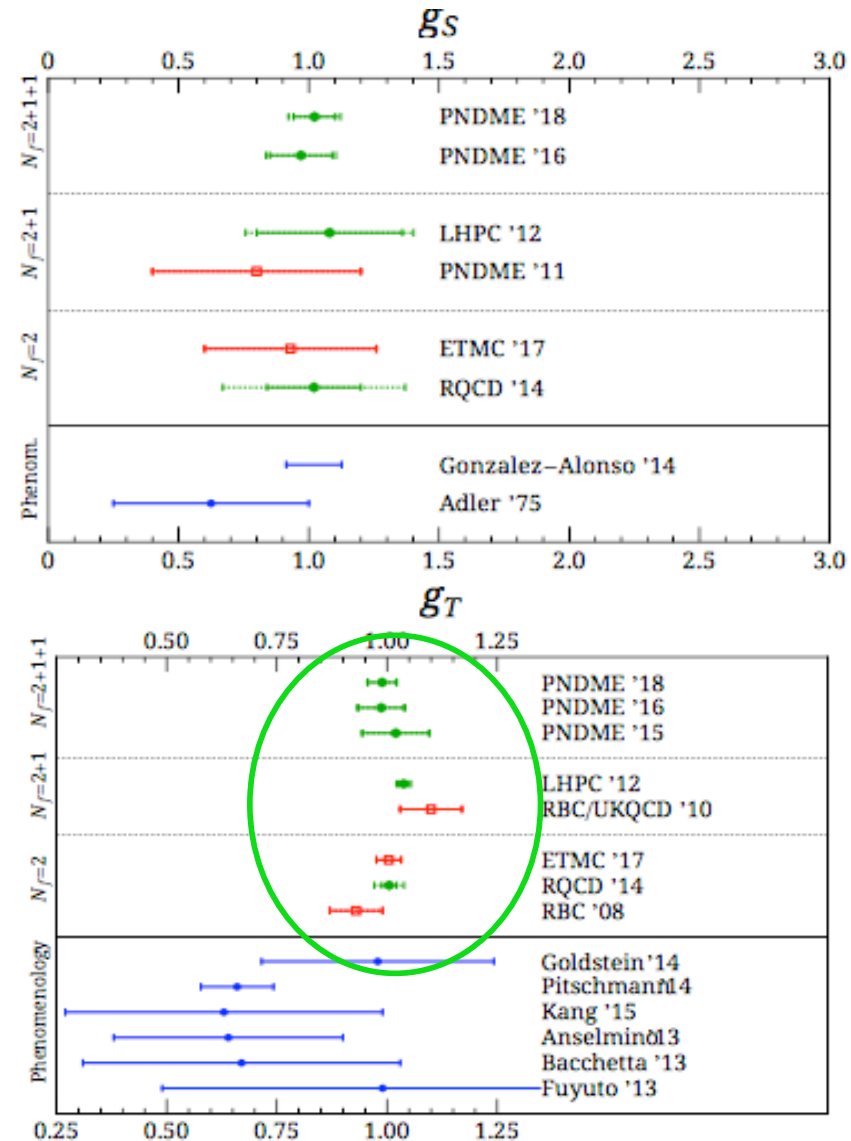
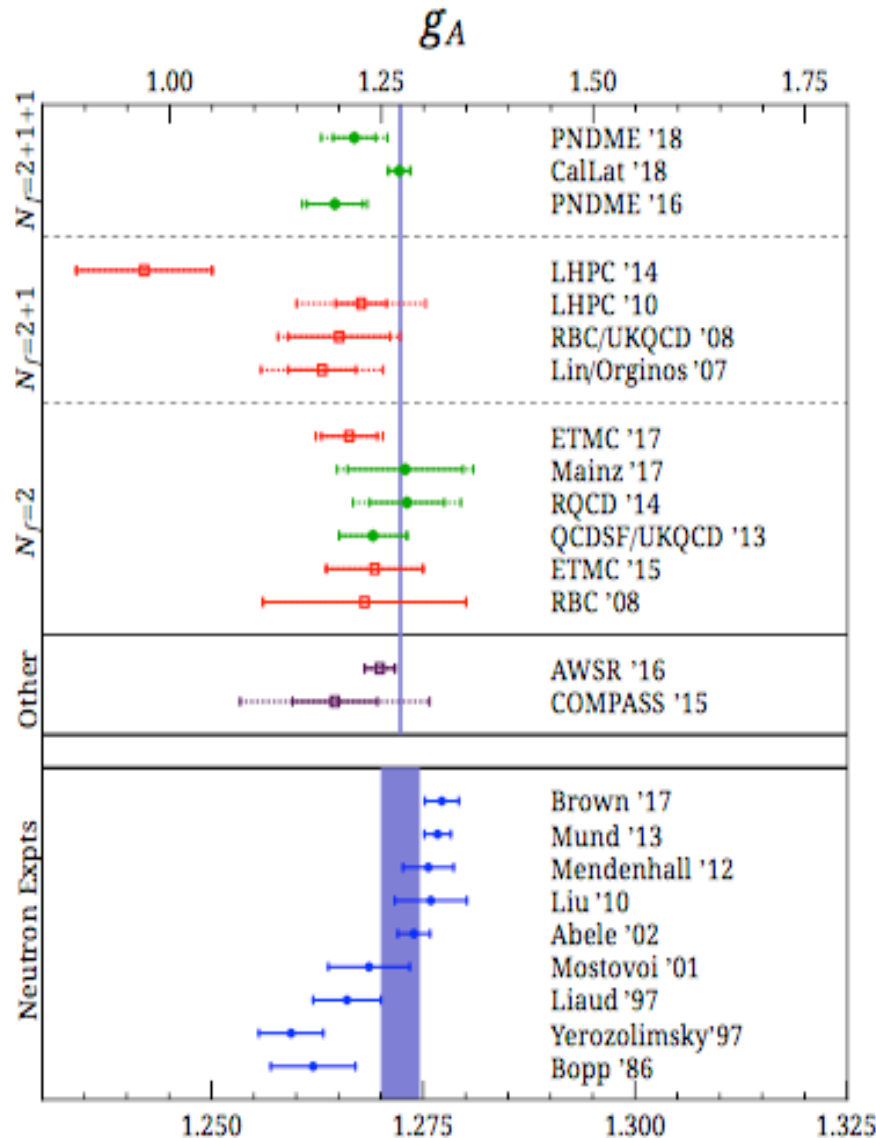
□ Axial charge:

	Lattice QCD	Global Fit
$g_A \equiv \langle 1 \rangle_{\Delta u^+ - \Delta d^+}$	$1.195(39) \ (N_f = 2 + 1 + 1)$	$1.275(12)$
	$1.279(50) \ (N_f = 2)$	

Moments of PDFs from lattice QCD

arXiv:1806.09006

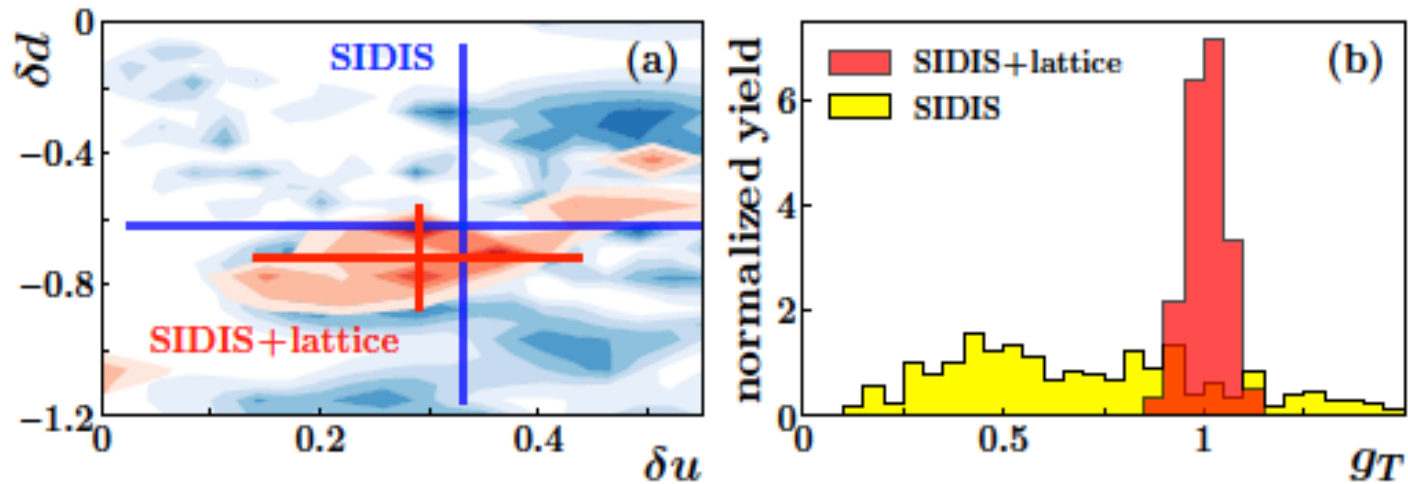
□ Isovector charges of nucleon: $g_A^{u-d}, g_S^{u-d}, g_T^{u-d}$



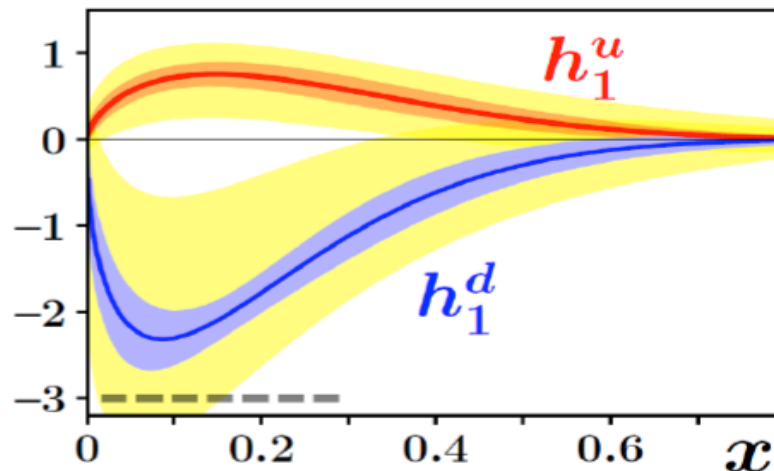
Lattice QCD helps QCD global analyses

□ Improved transversity distribution with LQCD g_T :

$$h_1(x, \mu^2) = \int \frac{dP^+ \xi^-}{2\pi} e^{-ixP^+ \xi^-} \langle P | \bar{\psi}(\xi^-) \frac{\gamma^+}{2P^+} \gamma_\perp \gamma_5 \exp \left\{ -ig \int_0^{\xi^-} d\eta^- A^+(\eta^-) \right\} \psi(0) | P \rangle$$



Much smaller
error band

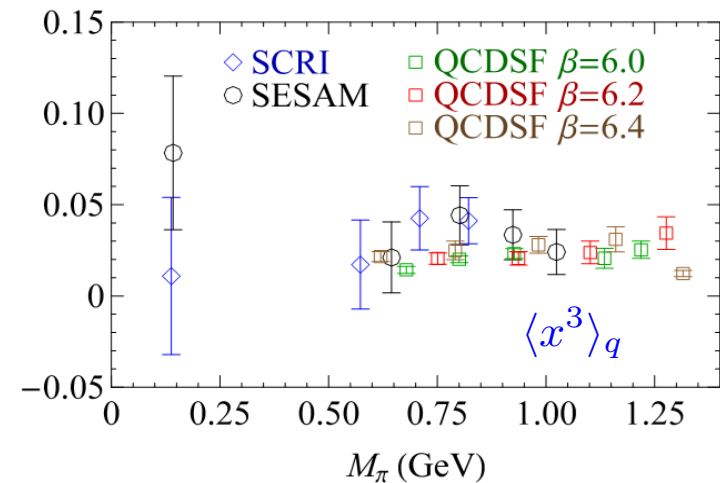
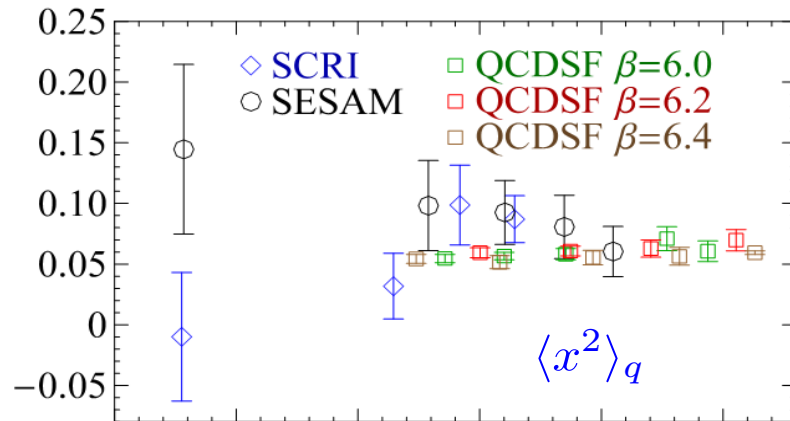


X-dependent PDFs from lattice QCD

Dolgov et al., hep-lat/0201021

Gockeler et al., hep-ph/0410187

□ Limited moments:



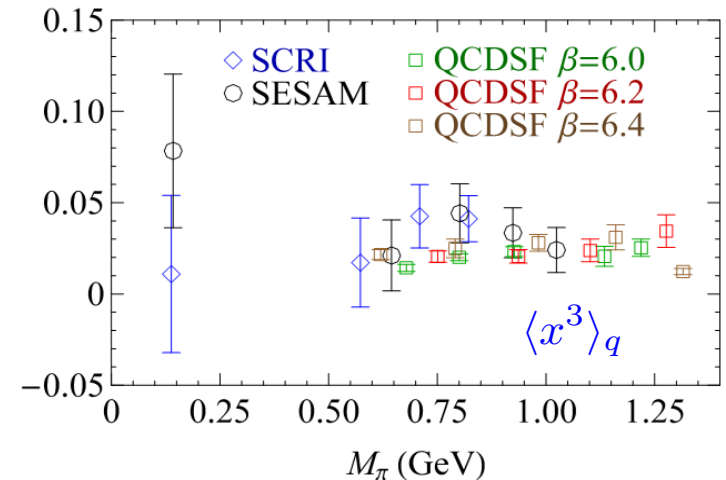
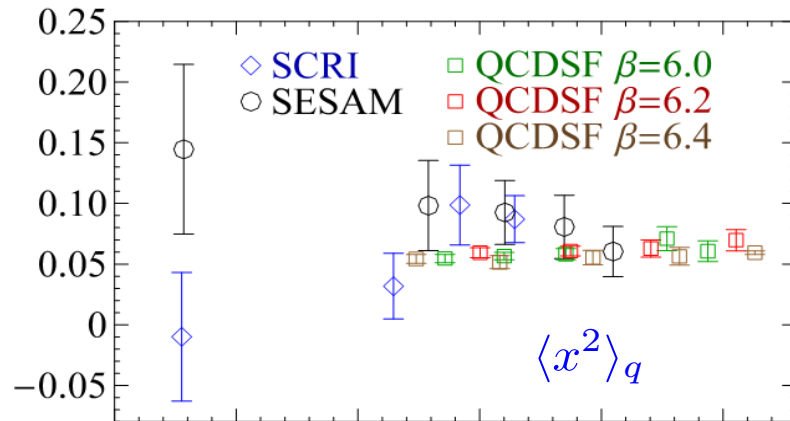
Limited moments – hard to get the full x-dependent distributions!

X-dependent PDFs from lattice QCD

□ Limited moments:

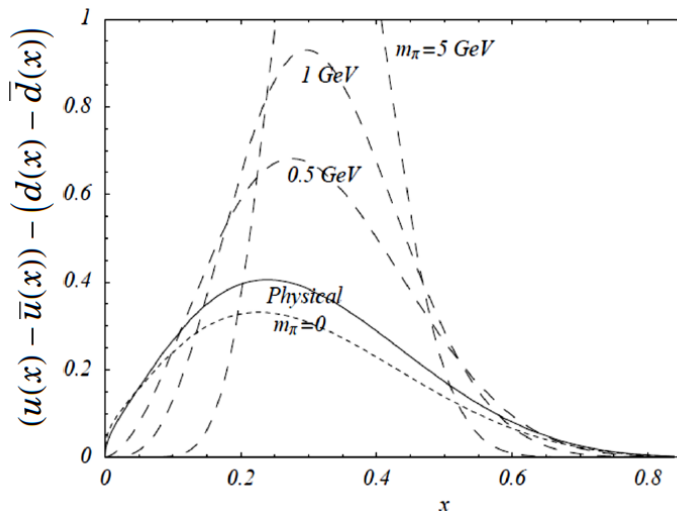
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Limited moments – hard to get the full x-dependent distributions!

□ Early efforts:



✧ Assume a smooth functional form:

$$xq(x) = a x^b (1 - x)^c (1 + \epsilon \sqrt{x} + \gamma x)$$

✧ Fix parameters with LQCD moments:

W. Dermold et al., Eur.Phys.J.direct C3 (2001) 1-15

Cannot distinguish valence quark contribution from sea quarks

New ideas

□ Quasi-PDFs:

Ji, arXiv:1305.1539

$$\tilde{q}(\tilde{x}, \mu_R^2, P_z) \equiv \int \frac{d\xi_z}{4\pi} e^{-i\tilde{x}P_z\xi_z} \langle P | \bar{\psi}(\frac{\xi_z}{2}) \gamma_z \exp \left\{ -ig \int_0^{\xi_z} d\eta_z A_z(\eta_z) \right\} \psi(\frac{-\xi_z}{2}) | P \rangle$$

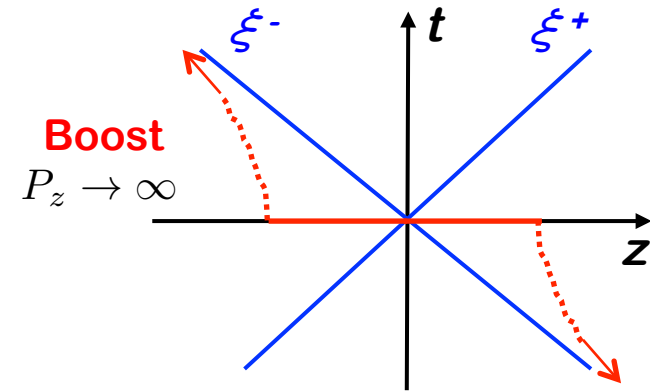
Idea:

Quasi-PDFs are not boost invariant

$$\tilde{q}(\tilde{x}, \mu_R^2, P_z) \longrightarrow q(x, \mu^2) \quad \text{when } P_z \rightarrow \infty$$

Note:

In Lattice QCD calculation, difficult to take $P_z \rightarrow \infty$ limit



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Proposed matching:

$$\tilde{q}(x, \mu^2, P_z) = \int_x^1 \frac{dy}{y} Z\left(\frac{x}{y}, \frac{\mu}{P_z}\right) q(y, \mu^2) + \mathcal{O}\left(\frac{\Lambda^2}{P_z^2}, \frac{M^2}{P_z^2}\right)$$

In terms of Large Momentum Effective Theory (LaMET)

Caution:

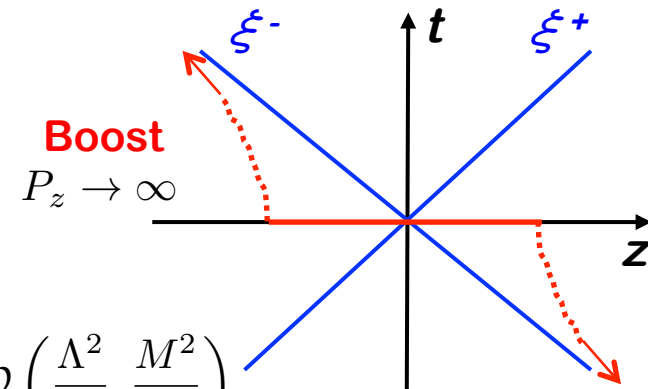
Power corrections could be of the form:

Braun, Vladimirov and Zhang, 1810.00048

$$\frac{\Lambda_{\text{QCD}}^2 R}{x^2(1-x)P_z^2}$$

Mixing with gluon and other flavor contribution beyond LO

Power UV divergence - μ_R does not obey DGALP



New ideas

□ Quasi-PDFs:

Ji, arXiv:1305.1539

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Progress:

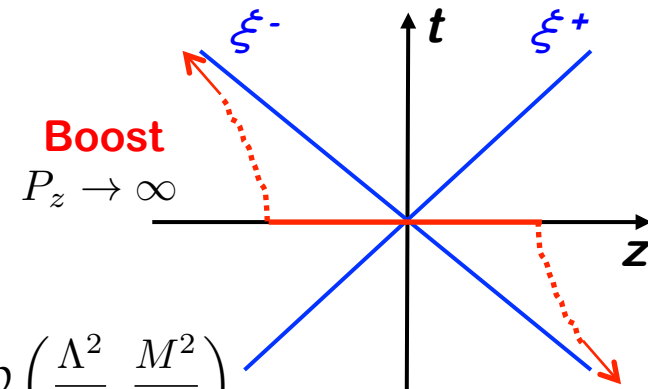
Quark: 1706.08962, 1707.03107, 1707.07152

Gluon: 1808.10824, 1809.01836

Power UV divergences are multiplicatively renormalizable

Renormalized Q-PDFs are collinearly factorizable to PDFs

01404.6860



New ideas

□ Pseudo-PDFs:

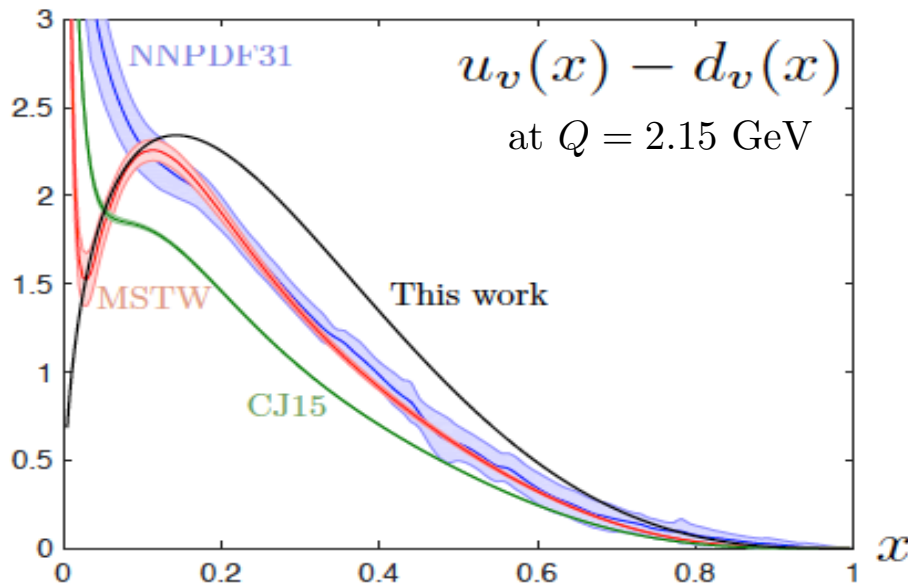
✧ Lattice calculation with $\alpha = 0$:

$$\begin{aligned}\mathcal{M}^\alpha(\nu = p \cdot \xi, \xi^2) &\equiv \langle p | \bar{\psi}(0) \gamma^\alpha \Phi_v(0, \xi, v \cdot A) \psi(\xi) | p \rangle \\ &\equiv 2p^\alpha \mathcal{M}_p(\nu, \xi^2) + \xi^\alpha (p^2/\nu) \mathcal{M}_\xi(\nu, \xi^2) \approx 2p^\alpha \mathcal{M}_p(\nu, \xi^2)\end{aligned}$$

$$\mathcal{P}(x, \xi^2) \equiv \int \frac{d\nu}{2\pi} e^{ix\nu} \mathcal{M}_{p=p^0}(\nu, \xi^2) / \mathcal{M}_{p=p^0}(0, \xi^2) \quad \leftarrow \text{Remove UV!}$$

✧ Model quasi-PDFs:

□ First numerical results:



Orginos, et al,
PRD96, 094503 (2017)

One-loop matching recently
Completed!

A. Radyushkin, arXiv:1801.02427

New ideas

Ma and Qiu, arXiv:1404.6860
arXiv:1709.03018

□ Good “Lattice cross sections”:

= Single hadron matrix element:

$$\sigma_n(\omega, \xi^2, P^2) = \langle P | T \{ \mathcal{O}_n(\xi) \} | P \rangle \quad \text{with } \omega \equiv P \cdot \xi, \quad \xi^2 \neq 0, \quad \text{and } \xi_0 = 0; \quad \text{and}$$

- 1) can be calculated in lattice QCD with precision,
has a well-defined continuum limit (UV+IR safe perturbatively), and
- 2) can be factorized into universal matrix elements of quarks and gluons
with controllable approximation

*Collaboration between lattice QCD
and perturbative QCD!*

New ideas

Ma and Qiu, arXiv:1404.6860
arXiv:1709.03018

□ Good “Lattice cross sections”:

= Single hadron matrix element:

$$\sigma_n(\omega, \xi^2, P^2) = \langle P | T \{ \mathcal{O}_n(\xi) \} | P \rangle \quad \text{with } \omega \equiv P \cdot \xi, \quad \xi^2 \neq 0, \quad \text{and } \xi_0 = 0; \quad \text{and}$$

- 1) can be calculated in lattice QCD with precision,
has a well-defined continuum limit (UV+IR safe perturbatively), and
- 2) can be factorized into universal matrix elements of quarks and gluons
with controllable approximation

*Collaboration between lattice QCD
and perturbative QCD!*

□ Current-current correlators:

$$\mathcal{O}_{j_1 j_2}(\xi) \equiv \xi^{d_{j_1} + d_{j_2} - 2} Z_{j_1}^{-1} Z_{j_2}^{-1} j_1(\xi) j_2(0)$$

with

d_j : Dimension of the current

Z_j : Renormalization constant of the current

Sample currents:

$$j_S(\xi) = \xi^2 Z_S^{-1} [\bar{\psi}_q \psi_q](\xi),$$

$$j_{V'}(\xi) = \xi Z_{V'}^{-1} [\bar{\psi}_q \gamma \cdot \xi \psi_{q'}](\xi),$$

$$j_V(\xi) = \xi Z_V^{-1} [\bar{\psi}_q \gamma \cdot \xi \psi_q](\xi),$$

$$j_G(\xi) = \xi^3 Z_G^{-1} [-\frac{1}{4} F_{\mu\nu}^c F_{\mu\nu}^c](\xi), \dots$$

□ Quasi- and pseudo-PDFs:

$$\mathcal{O}_q(\xi) = Z_q^{-1} (\xi^2) \bar{\psi}_q(\xi) \gamma \cdot \xi \Phi(\xi, 0) \psi_q(0)$$

$$\Phi(\xi, 0) = \mathcal{P} e^{-ig \int_0^1 \xi \cdot A(\lambda \xi) d\lambda}$$

Exploratory effort and its success

❑ Two lattice collaborations did pioneering work for quark-qPDFs:

Both LP3 and ETMC obtained their results at physical pion mass!

✧ Calculate quasi-quark distribution:

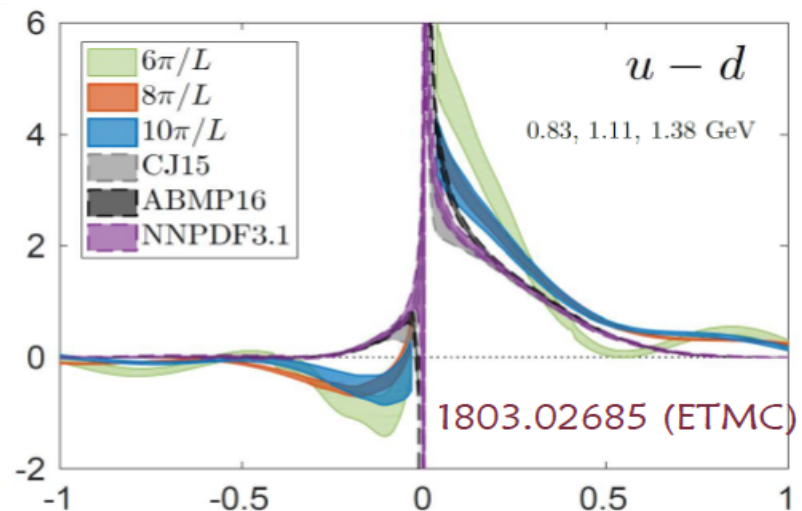
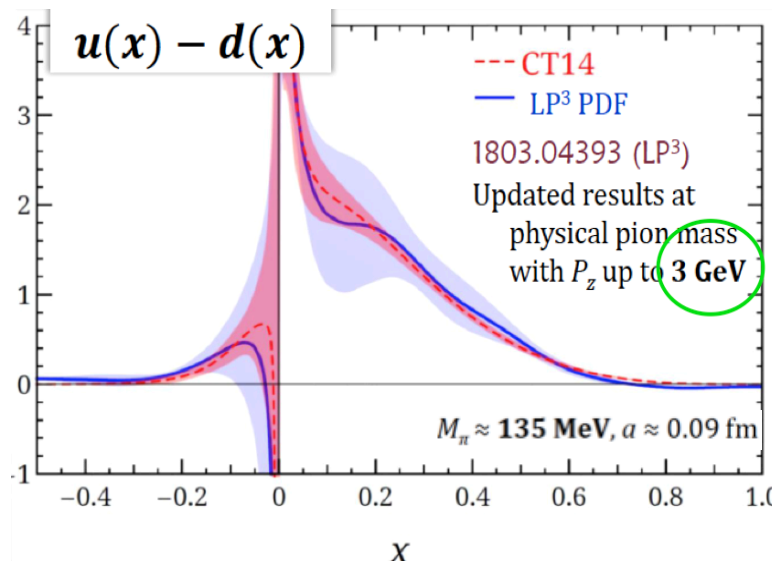
$$\tilde{q}(x, \mu, |P_z) = \int \frac{dz}{4\pi} e^{-izk_z} \left\langle P \left| \bar{\psi}(z) \Gamma \exp\left(-ig \int_0^z dz' A_z(z')\right) \psi(0) \right| P \right\rangle$$

✧ Extract quark distribution:

$$\tilde{q}(x, \mu, P_z) = \int_{-\infty}^{\infty} \frac{dy}{|y|} Z\left(\frac{x}{y}, \frac{\mu}{P_z}\right) q(y, \mu) + \mathcal{O}(M_N^2/P_z^2) + (\Lambda_{\text{QCD}}^2/P_z^2)$$

✧ The state of art:

Inverse to get the PDFs

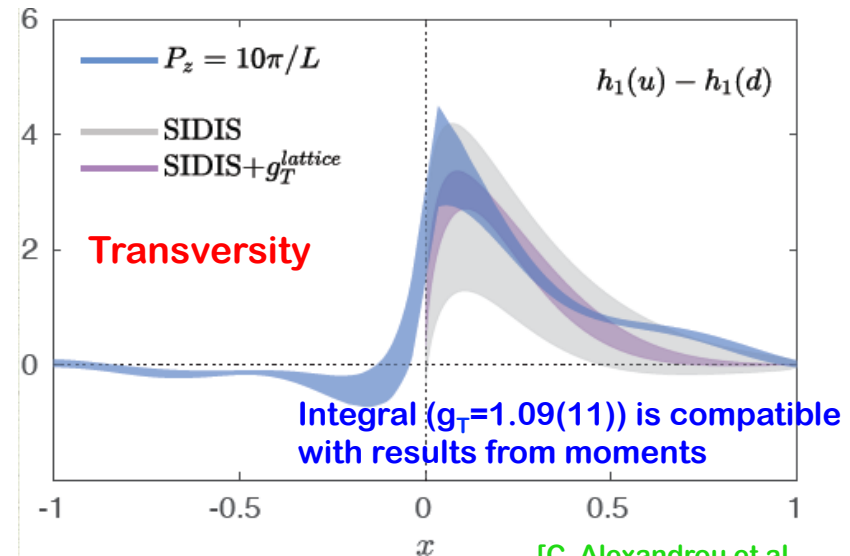
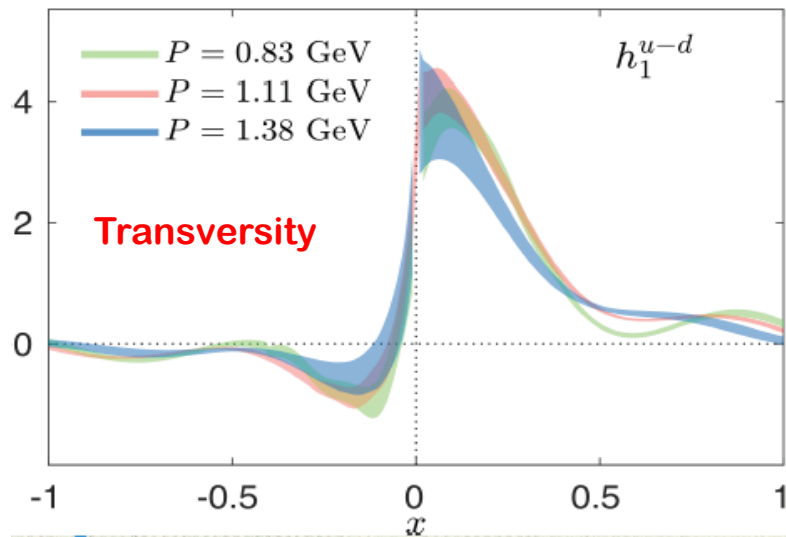
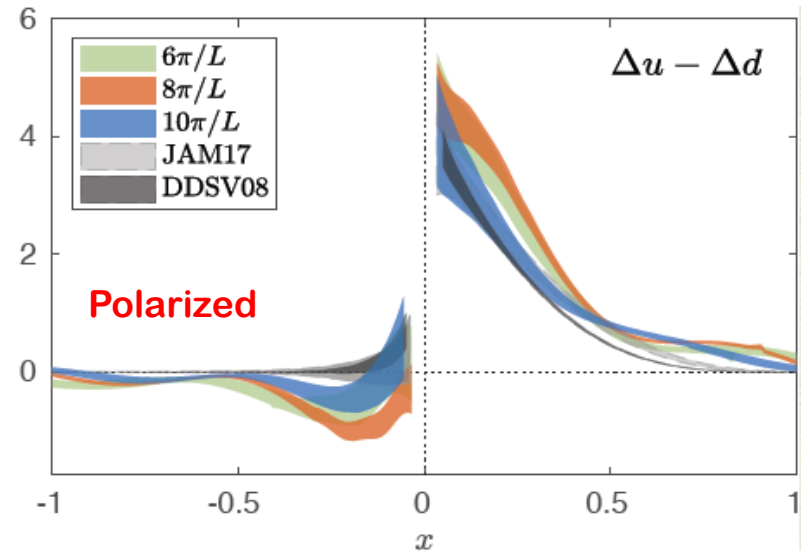
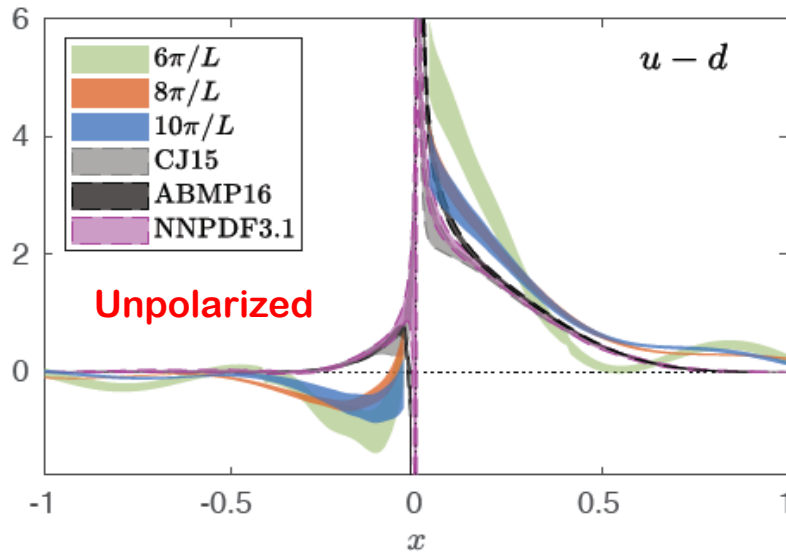


With one-loop perturbative matching and target mass correction

Exploratory effort and its success

□ Light-cone PDFs - ETMC:

M. Constantinou at CTEQ meeting



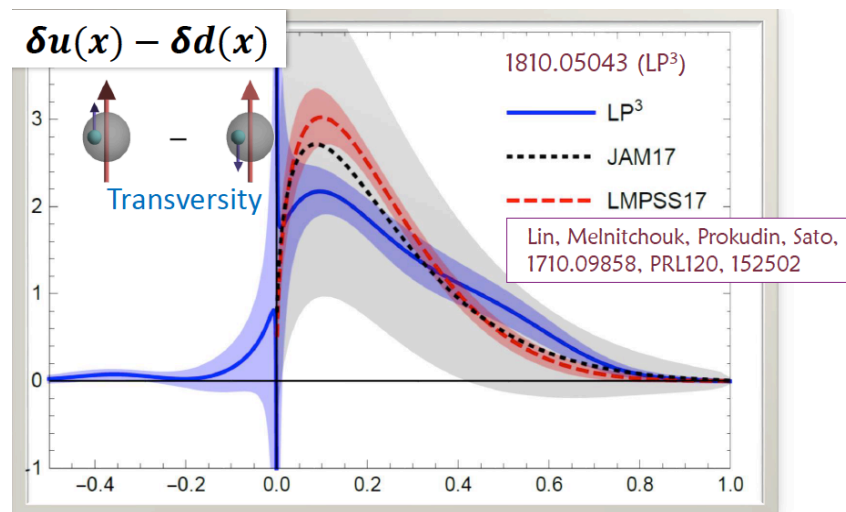
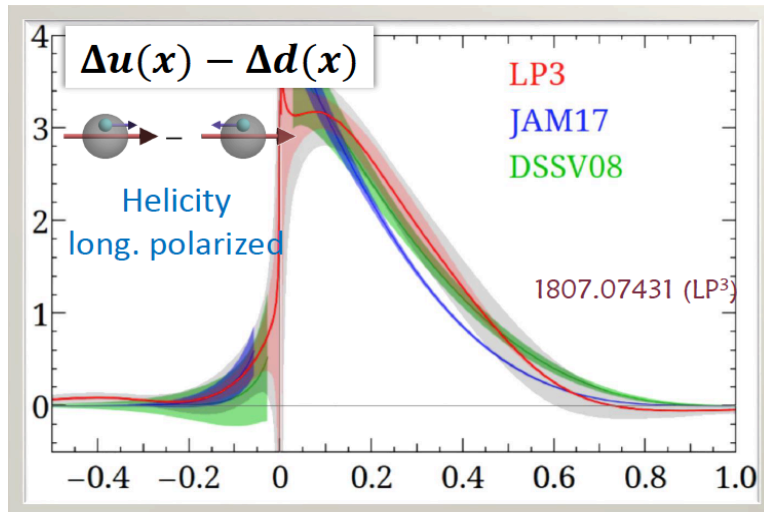
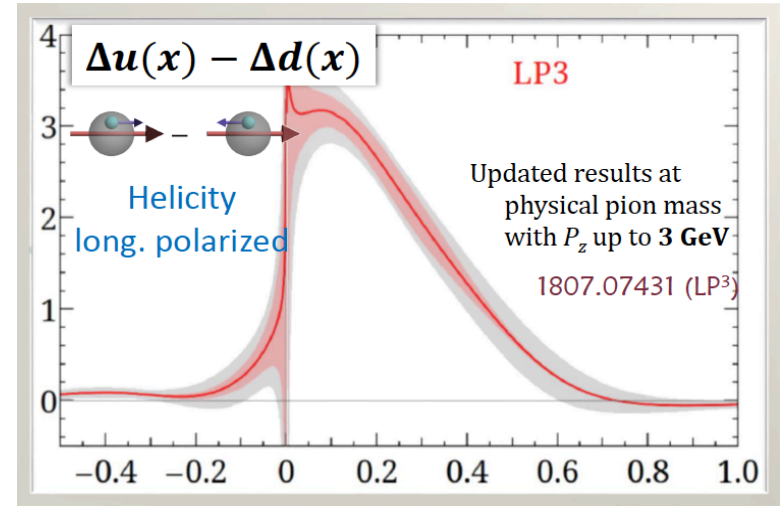
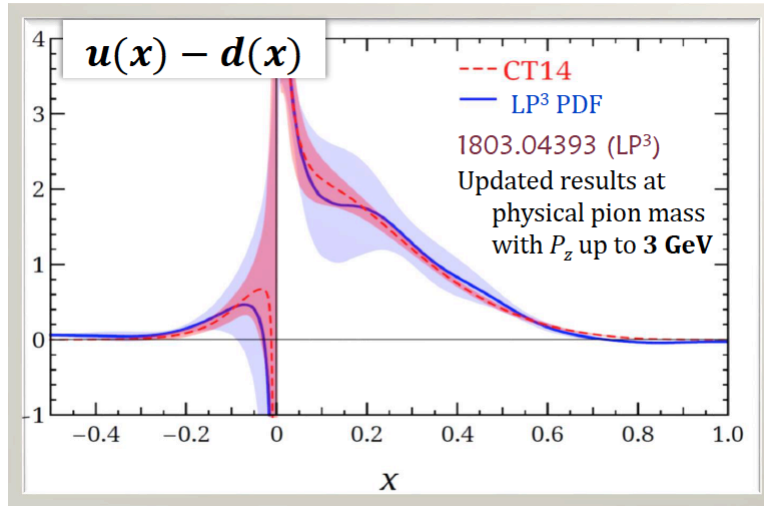
All with one-loop perturbative matching!

[C. Alexandrou et al., PRD95, 114514 (2017)]

Exploratory effort and its success

□ Light-cone PDFs – LP3:

H.W. Lin at CTEQ meeting



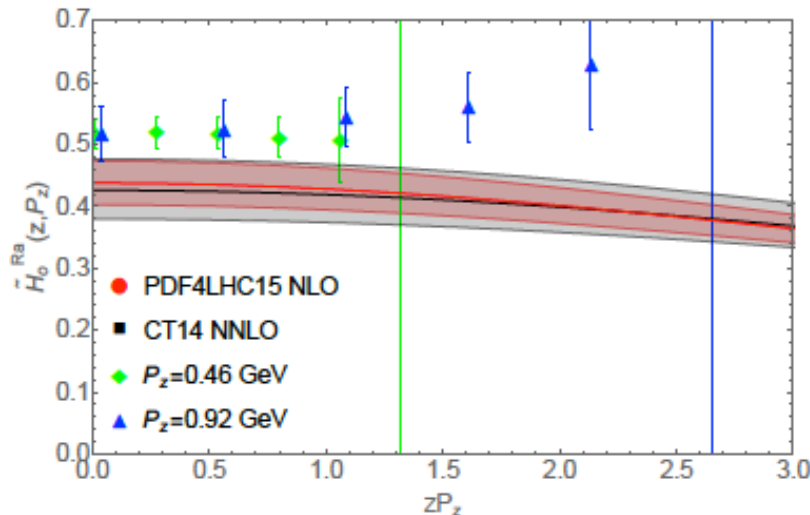
All with one-loop perturbative matching!

Exploratory effort and its success

□ First glimpse into gluon Q-PDFs – LP3:

H.W. Lin at CTEQ meeting
1808.02077

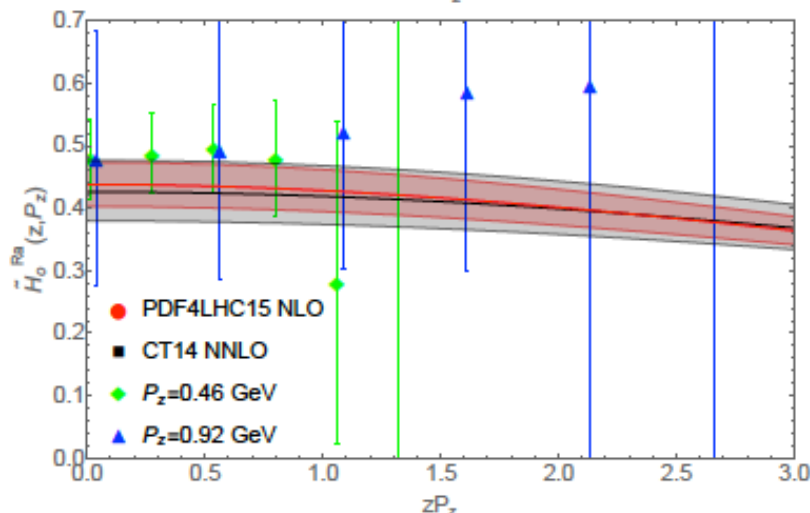
Lattice details: overlap/2+1DWF, 0.16fm, 340-MeV sea pion mass



Like quark Q-PDFs, gluon Q-PDFs have power UV divergence, and could mix with others under renormalization, ...

UV divergence are multiplicatively
Renormalizable to all orders in pQCD

Zhang et al, arXiv:1808.10824
Li et al, arXiv:1809.01836

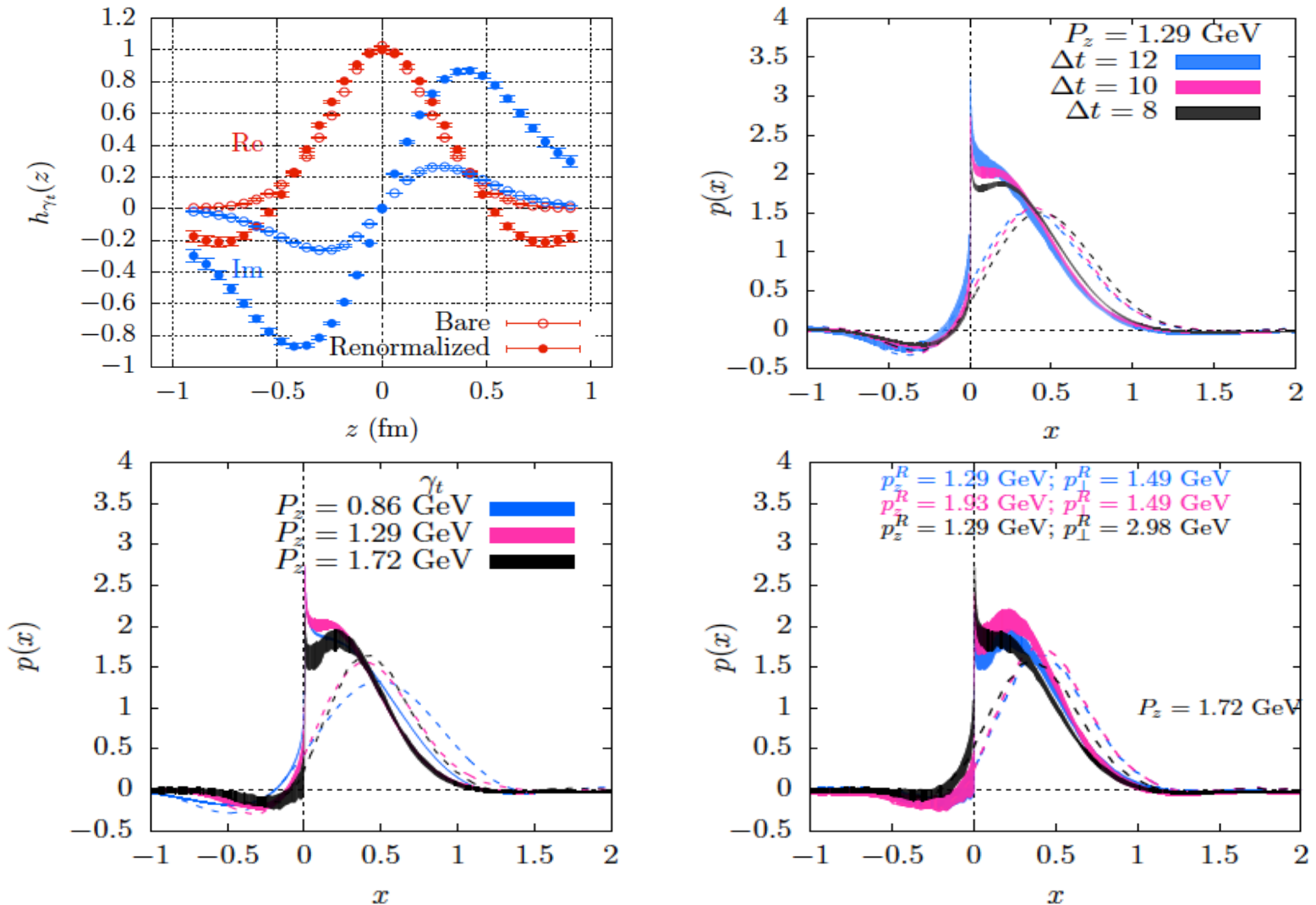


$$\tilde{H}_0^{Ra}(z, P_z, \mu) = \frac{\tilde{H}_0^{\overline{\text{MS}}}(0, 0, \mu)}{\tilde{H}_0(z, 0)} \tilde{H}_0(z, P_z)$$

Exploratory effort and its success

□ Pion qPDFs & PDFs:

N. Karthik at Lattice2018
BNL, SB, UConn, ...

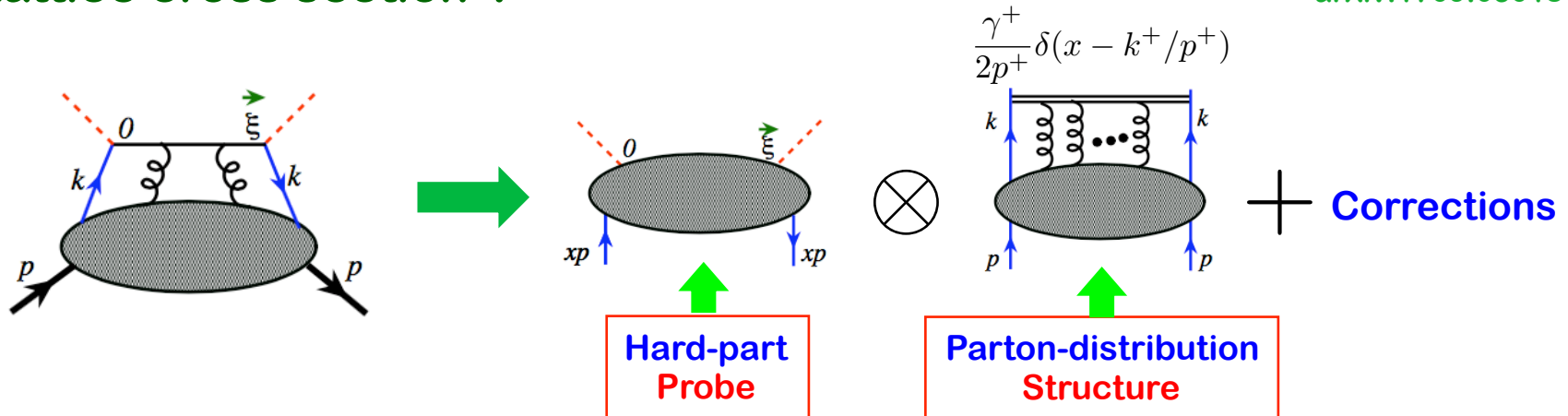


All with one-loop perturbative matching!

Exploratory effort and its success

❑ “Lattice cross section”:

Ma and Qiu, arXiv:1404.6860
arXiv:1709.03018



❑ QCD Global analysis:

Need data of “many” good lattice cross sections to be able to extract the x , Q , flavor dependence of the structure, ...

❑ Complementarity and advantages:

- ✧ Complementary to existing approaches for extracting PDFs,
- ✧ Quasi-PDFs and pseudo-PDFs are special cases,
- ✧ Have tremendous potentials:

Neutron PDFs, ... (no free neutron target!)

Meson PDFs, such as pion, ...

More direct access to gluons – gluonic current, quark flavor, ...

Exploratory effort and its success

□ Pion/Keon PDFs:

JLab lattice group
Ma, Qiu, arXiv:1709.03018
PRL (2018)

– using a vector-axial-vector correlation as an example

✧ Parity-Time-reversal invariance:

$$\sigma_{ij}^{\mu\nu}(\xi, p) = \xi^4 \langle \pi(p) | \mathcal{J}_i^\mu(\xi/2) \mathcal{J}_j^\nu(-\xi/2) | \pi(p) \rangle$$

$$\frac{1}{2} [\sigma_{VA}^{\mu\nu}(\xi, p) + \sigma_{AV}^{\mu\nu}(\xi, p)] \equiv \epsilon^{\mu\nu\alpha\beta} \xi_\alpha p_\beta T_1(\omega, \xi^2) + (p^\mu \xi^\nu - \xi^\mu p^\nu) T_2(\omega, \xi^2)$$

✧ Collinear factorization:

$$T_i(\omega, \xi^2) = \sum_{a=q, \bar{q}, g} \int_0^1 \frac{dx}{x} f_a(x, \mu^2) C_i^a(x\omega, \xi^2, \mu^2) + \mathcal{O}(\xi^2 \Lambda_{\text{QCD}}^2)$$

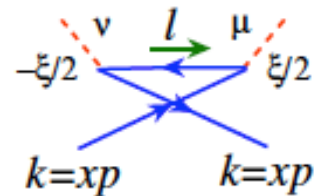
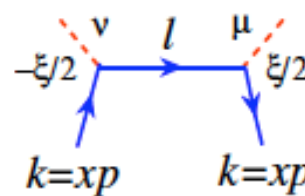
✧ Lowest order coefficient functions:

$$C_1^{q(0)}(x\omega, \xi^2) = T_1^{q(0)}(x\omega, \xi^2) = \frac{2x}{\pi^2} \cos(x\omega)$$

$$C_2^{q(0)}(x\omega, \xi^2) = 0$$

$$\tilde{T}_1(\tilde{x}, \xi^2) \equiv \int \frac{d\omega}{2\pi} e^{-i\tilde{x}\omega} T_1(\omega, \xi^2)$$

$$\approx \frac{1}{\pi^2} \{ q(\tilde{x}) - \bar{q}(\tilde{x}) \} = \frac{1}{\pi^2} q_v(\tilde{x})$$

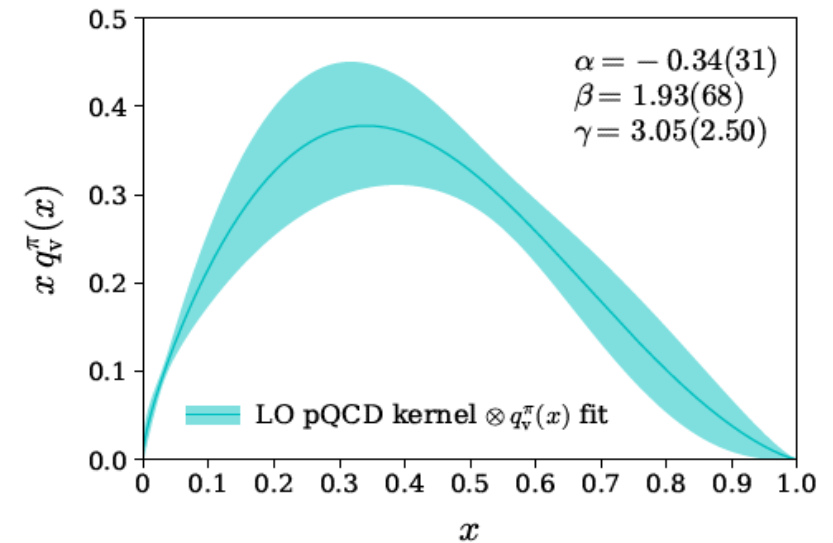
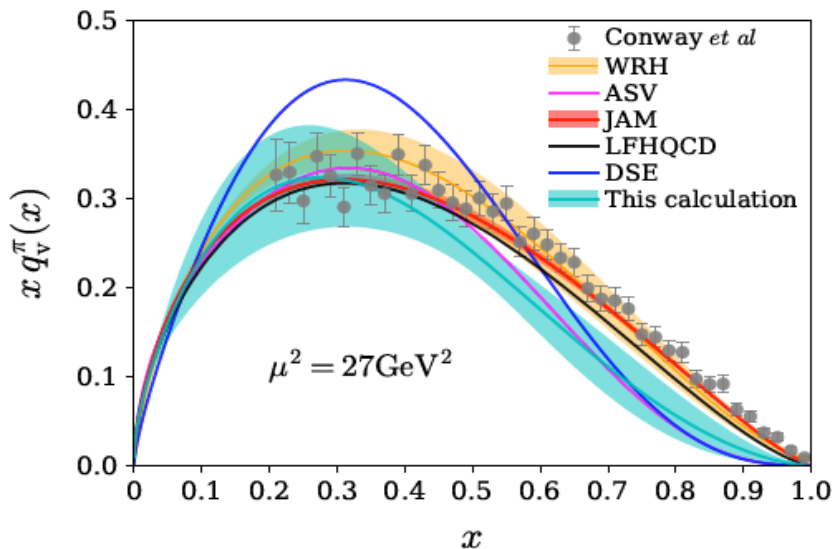
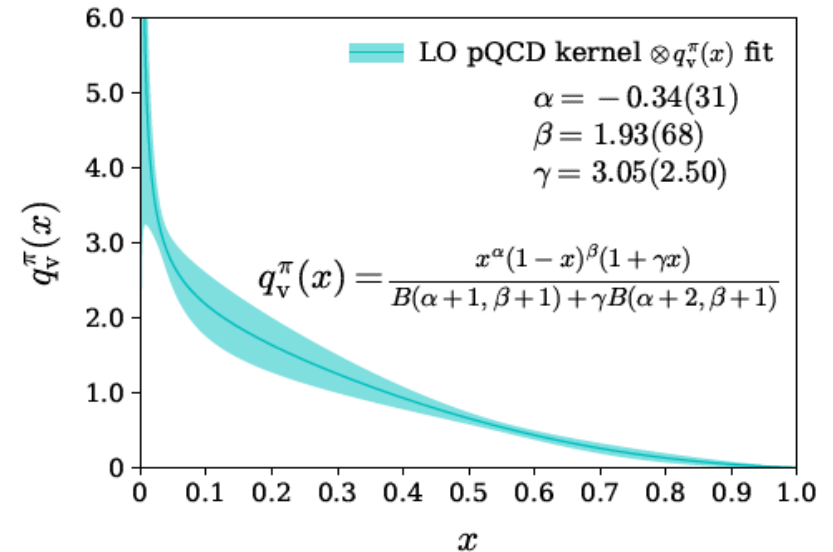
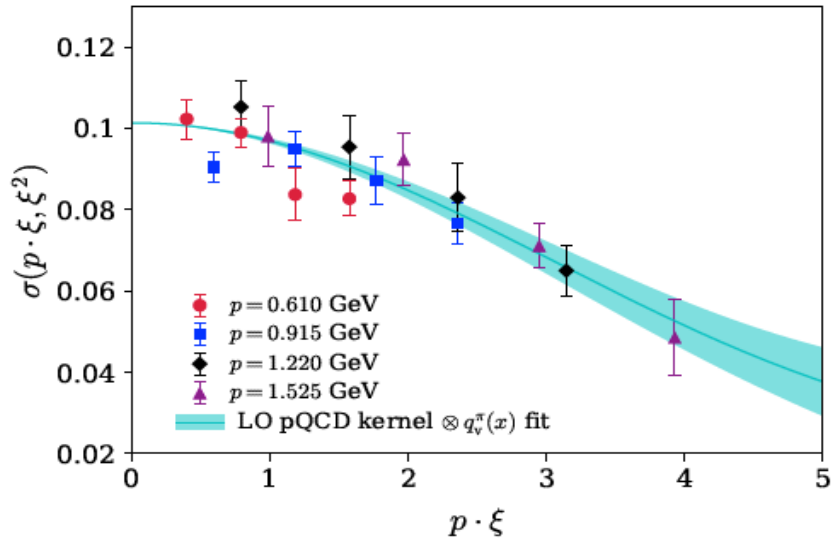


Lattice QCD calculation → pion valence quark distribution!

Exploratory effort and its success

JLab lattice group
arXiv:1901.03921

□ Pion valence PDFs:



Tree-level (LO) hard coefficient only!

Summary and outlook

□ Lattice QCD calculation will provide rich information on hadron's partonic structure:

- ✧ Quasi-PDFs approach
- ✧ Pseudo-PDFs approach
- ✧ “Lattice cross section” approach

Community White Paper
arXiv:1711.07916

Key: Controllable matching to PDFs, GPDs, TMDs, ...

Summary and outlook

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Community White Paper
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Key: Controllable matching to PDFs, GPDs, TMDs, ...

- Results of initial exploratory calculations are encouraging:

- ✧ Good for relatively large x partonic structure – computing power
- ✧ Complementary to experimental measurements

$$\tilde{q}(\tilde{x}, \mu_R^2, P_z) \equiv \int \frac{d\xi_z}{4\pi} e^{-i\tilde{x}P_z\xi_z} \langle P | \bar{\psi}(\frac{\xi_z}{2}) \gamma_z \exp \left\{ -ig \int_0^{\xi_z} d\eta_z A_z(\eta_z) \right\} \psi(\frac{-\xi_z}{2}) | P \rangle$$

With $P_z \sim \text{GeV}$, $\xi_z \sim \text{fm}$, it is hard for $x < 0.1$ while keeping the exponential to be larger than 1!

- Like its success in calculating the “baryon interaction” and “hadron spectroscopy”, lattice QCD can be used to study hadron structure, while more works are needed!

Thank you!